

MISOLFA solar monitor for the ground PICARD program

T. Corbard^{1,*}, A. Irbah², P. Assus¹, C. Dufour², M. Fodil³, F. Morand¹, C. Renaud¹, and E. Simon¹

¹ Université de Nice Sophia Antipolis, CNRS, Observatoire de la Côte d’Azur, BP 4229 06304 Nice Cedex 4, France

² Université Versailles St-Quentin, CNRS/INSU, LATMOS-IPSL, Guyancourt, FRANCE

³ CRAAG, Observatoire d’Alger, BP 63 Bouzaréah 16340 Alger, Algérie

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Developed at the Observatoire de la Côte d’Azur (OCA) within the framework of the PICARD space mission (Thuillier et al., 2006) and with support from the french spatial agency (CNES), MISOLFA (Moniteur d’Images Solaires Franco-Algérien) is a new generation of daytime turbulence monitor. Its objective is to measure both the spatial and temporal turbulence parameters in order to quantify their effects on the solar diameter measurements that will be made from ground using the qualification model of the SODISM (Solar Diameter Imager and Surface Mapper) instrument onboard PICARD. The comparison of simultaneous images from ground and space should allow us, with the help of the solar monitor, to find the best procedure possible to measure solar diameter variations from ground on the long term. MISOLFA is now installed at the Calern facility of OCA and PICARD is scheduled to be launched in 2010. We present here the principles of the instrument and the first results obtained on the characteristics of the turbulence observed at Calern observatory using this monitor while waiting for the launch of the space mission.

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1 Introduction: why a new Solar monitor ?

Diameter measurement analysis revealed a dependence with the seeing condition as represented by Fried’s parameter r_0 (Irbah et al., 1994). Numerical simulations were developed in order to better understand atmospheric effects on diameter measurements. The error decreases with the seeing but it is also strongly conditioned by turbulence coherence times (Fig.2, Lakhali et al., 1999). The error also shows weak dependence on the outer scale L_0 for a small aperture telescope. Existing solar monitors such as Solar Differential Motion Monitors (S-DIMM) or arrays of scintillometers (SHABAR) are able to provide useful information on the spatial scales of the turbulence and are commonly associated for site testing (e.g. Beckers 2001). However our goal here is to obtain the accurate turbulence equivalent PSF needed to properly interpret the ground based radius measurements. For this, estimates of the characteristic temporal scales are also needed and it was realized that a single instrument could provide both the spatial and temporal turbulence scales by analysing Angle of Arrival (AA) fluctuations simultaneously in its image and pupil ways.

2 Atmospheric parameters measured

An illustration of a multi-layer atmosphere model is given on Fig. 3. MISOLFA will allow us to obtain information on



Fig. 1 Picture of MISOLFA monitor and the mount that will receive the qualification model of the PICARD/SODISM telescope at the Calern facility of OCA.

all the following parameters using a Von Kármán turbulence model:

- The Atmospheric structure constant of the air refractive index fluctuations $C_n^2(h)$.
- Fried’s parameter r_0 which is the diameter of the coherence zone of the degraded wavefront. It corresponds also to the image resolution obtained with the telescope of diameter r_0 placed outside the atmosphere.
- The spatial coherence outer scale L_0 which defines the maximal size of wavefront perturbations remaining coherent. It traduces the low frequency evolution of the wavefront.
- The isoplanatic patch θ_0 which is the angle where AA or speckles remain correlated.

* Corresponding author: e-mail: Thierry.Corbard@oca.eu

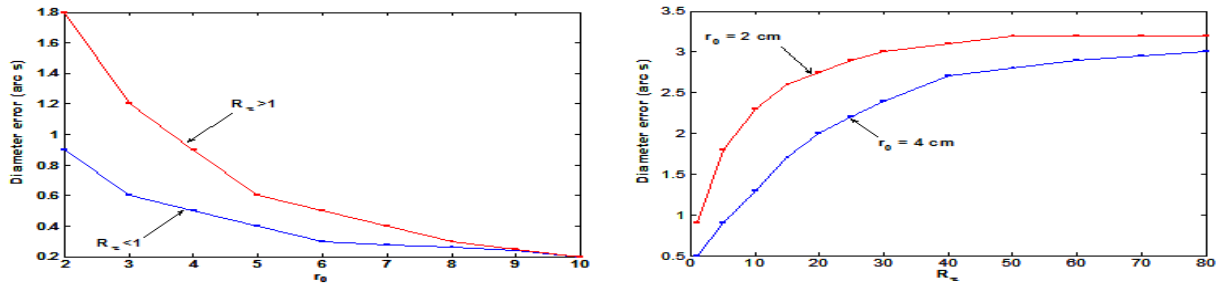


Fig. 2 Estimated errors on diameter estimates as a function of the Fried's parameter r_0 for two values of R_τ (ratio between the exposure time and the correlation time) (left) and as a function of R_τ for two values of r_0 (right) (Lakhal et al., 1999)

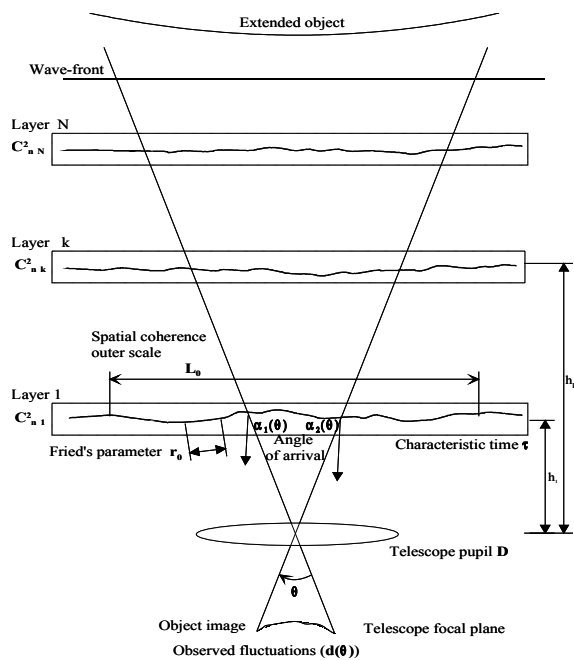


Fig. 3 Illustration of a multi-layers atmosphere with Von Kármán turbulence model (Irbah et al., 2003).

- The correlation time τ_0 which is the time during which the atmosphere may be considered as frizzed for the considered structures (AA, speckles) i.e. the time during which they keep their coherence.

3 Measurement principles

The MISOLFA principle is based on the statistical analysis of AA fluctuations, which are fluctuations at each point of the normal of the perturbed wavefronts. The AA fluctuations can directly be observed in the image plane (case of Shack-Hartmann's sensors used in adaptive optics) but also in the pupil plane if the observed source present an intensity distribution with a strong discontinuity like the Solar limb. The intensity fluctuations observed in the pupil image are,

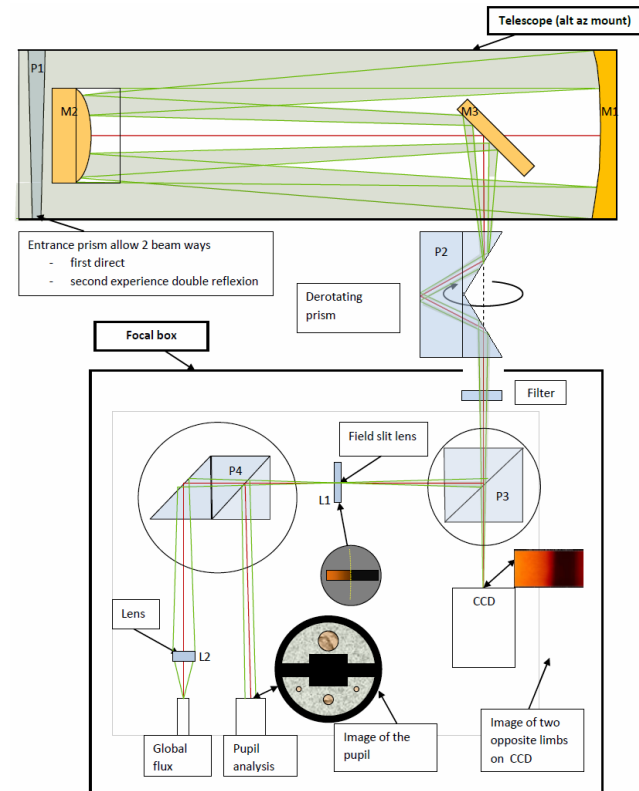


Fig. 4 MISOLFA optical design. See text for details.

at first order, proportional to AA fluctuations (Borgnino & Martin 1977, Borgnino 1978).

4 MISOLFA Instrument design

Figure 4 illustrates the optical design of MISOLFA. It is made of a Cassegrainian telescope of 25 cm aperture equipped with an alt-azimuth mount to which is associated a derotating system. A flat glass blade with non parallel faces allows the formation of separate images of the sun (Fig. 5). In the focal box, a first beam (image way) is imaging two opposite parts of solar limb on a 480x640 CCD. A typical sequence is 1.5 ms long with 32 images per second and a resolution

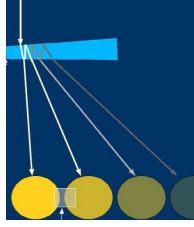


Fig. 5 Illustration of the formation of several solar images by the entrance prismatic blade. The CCD camera is positioned across the two first solar images.

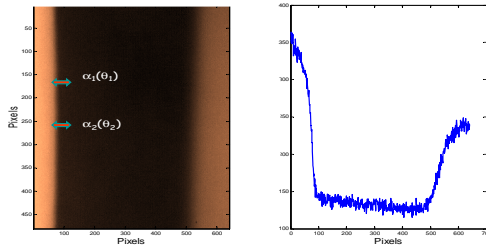


Fig. 6 Image sample from the MISOLFA monitor (left) and horizontal cut (right). Arrows illustrate the motions used to compute the transverse covariance between to angular positions (see text).

of 0.2 arcseconds per pixel. A second beam (pupil way) is imaging the pupil after passing through a spatial filter (slit) in the image of a part of solar limb. In the pupil image 4 optical fibers (with diameters 2 mm, 1 mm and 2x0.5 mm) guide the light to 4 diodes. In addition, in the image of the slit a 2 mm fiber measure global flux passing through the slit.

5 Analysis Method

In the image way we estimate r_0 , L_0 and θ_0 from direct measurement of AA fluctuations on the limb images:

- r_0 is directly related to the variance σ_s^2 of the mean AA fluctuations by:

$$r_0 = 8.2510^5 D^{-1/5} \lambda^{6/5} (\sigma_s^2)^{-3/5}$$

where D is the pupil diameter and λ the wavelength.

- The AA transverse covariance is computed from the observed limb fluctuations by (Fig. 6):

$$C_\alpha(\theta) = \langle \alpha_1(\theta_1) \alpha_2(\theta_2) \rangle$$

L_0 and $\theta_0 = r_0/h$ are then obtained by adjusting the observed covariance to the theoretical one given, in the case of the one-layer model and Von Kármán theory, by:

$$C_\alpha(\theta) = 0.0716 \lambda^2 r_0^{-5/3} \int_0^\infty f^3 \left(f + \frac{1}{L_0^2} \right)^{-\frac{11}{6}} \dots \\ \dots [J_0(2\pi f \theta h) + J_2(2\pi f \theta h)] \left[\frac{2J_1(\pi D f)}{\pi D f} \right]^2 df$$

where h is the altitude of the equivalent impulse layer giving at ground level the same optical effects than the whole turbulent terrestrial atmosphere.

- The turbulence profile $C_n^2(h)$ is then obtained by inverting (Bouzig et al., 2000) the structure function which, in a multi-layer model, is given by:

$$D_\alpha(\theta) = 2 [C_\alpha(0) - C_\alpha(\theta)] \\ = 2.4 \int_0^\infty C_n^2(h) K(\theta, h, D, L_0) dh$$

with:

$$K(\theta, h, D, L_0) = \int_0^\infty f^3 \left(f + \frac{1}{L_0^2} \right)^{-\frac{11}{6}} \dots \\ \dots [1 - J_0(2\pi f \theta h) - J_2(2\pi f \theta h)] \left[\frac{2J_1(\pi D f)}{\pi D f} \right]^2 df$$

In the pupil plane, we perform a spatio-temporal analysis of AA fluctuations (Berdja, 2007). The photodiode detectors used represent various collecting surfaces of the pupil-plane. They allow fast recordings (5KHz) and thus an accurate intensity signal time sampling. Such measurement time series will be used to study temporal properties of the diurnal turbulence and continuously estimate the correlation time τ_0 of AA-fluctuations from the estimated temporal covariance assuming the Von Kármán model and Taylor's hypothesis of frozen turbulence. Such as for other seeing monitors like the well-known S-DIMM or GSM, dual measurements from photodiode pairs on the same collecting surface will allow also estimate the Fried parameter r_0 . The needed transverse temporal covariance is obtained by doing $\theta h = v\tau$ in the previous formula, v denoting an equivalent speed (weighted average of the different layer speeds) and τ a temporal shift (Conan et al., 2000). The goal of using photodiode detectors with different collecting surfaces is to estimate the spatial coherence outer scale L_0 from AA-fluctuation statistics.

6 Observations and first results

Sequences of images are recorded at Calern observatory since June 2009. A typical sequence consists in a set of about 2500 images with an angular resolution of 0.2 arcseconds per pixel recorded at a rate of 32 images per second. The exposure time of each image is 1 ms. The mean limb position is evaluated for each image in order to correct in real time the guiding errors. The limb profile of each image is first extracted using wavelet based denoising and inflexion point identification (Irbah et al. 1999) as illustrated Fig. 7. Then, from the temporal evolution of all the points of the extracted limb profiles (Fig. 8), the variance of the mean AA fluctuations and the Fried's parameter are estimated. When estimating the Fried's parameter from increasing time intervals, the value decreases while integrating clearly different turbulence regimes (Fig. 9). Figure 10 shows Fried's parameter r_0 estimated on a continuous set of two seconds

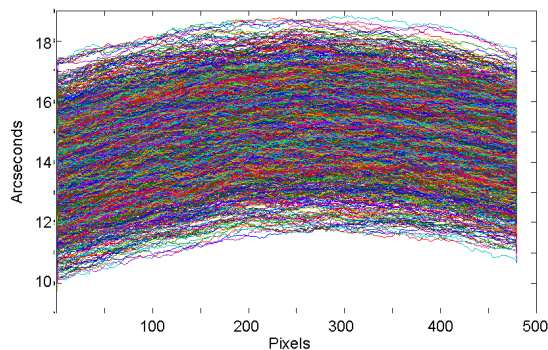


Fig. 7 The limb profile of each image is extracted using wavelet based denoising and inflexion point identification (Irbah et al. 1999)

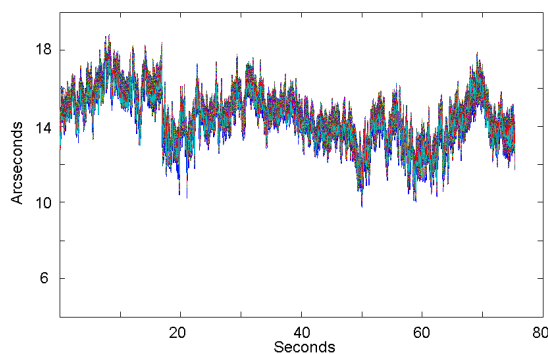


Fig. 8 Temporal evolution of all the points of the extracted limb profiles.

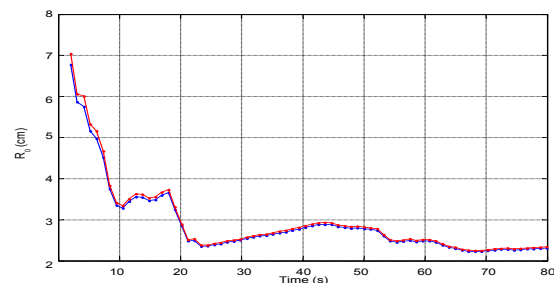


Fig. 9 Fried's parameter r_0 estimated from increasing time intervals.

intervals which correspond to the exposure time of the full disk images that will be recorded by the SODISM instrument. The full time sequence is therefore cutted in 40 intervals of two seconds. On these time intervals r_0 can reach up to 14 cm while over the whole 75 s the value would be about 2.6 cm (Fig. 9).

7 Conclusion

MISOLFA is a new type of solar monitor based on the analysis of AA statistic that we record both in the image and pupil plane. It will allow us to get a full characterization

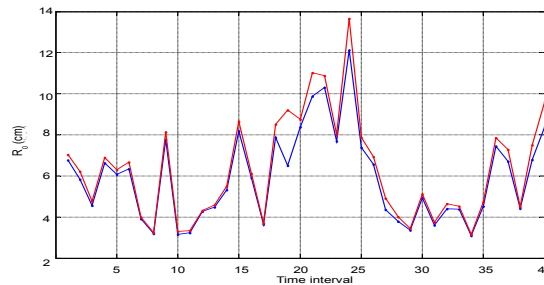


Fig. 10 Fried's parameter r_0 estimated on two seconds intervals.

of the spatio-temporal parameters of the turbulence needed to identify the appropriate atmospheric model and build the turbulence + instrument equivalent PSF. The first sequences of images obtained were analyzed to extract the Fried's parameter showing that good condition can eventually be reached at Calern observatory on some intervals for the 2 s exposure time of SODISM. The continuous record of the r_0 values will help us in the interpretation of ground SODISM images while the PICARD satellite is operating. The procedure to extract all the other spatial turbulence parameters have been tested and are now being implemented. The pupil way is still not operating properly and lots of effort are currently made to increase the signal to noise ratio by improving the electronic components.

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